

Correlation Functions in Classical Statistical Physics

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$$U_n(f_1, \dots, f_n) = \frac{\partial^n}{\partial t_n \dots \partial t_1} \log \left(\mu(e^{\sum_{i=1}^n t_i f_i}) \right) \Big|_{t_1=\dots=t_n=0} .$$

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► For the sake of concreteness, consider the Ising model on $\mathbb{Z}^d$ given by the (formal) Hamiltonian:

$$H = - \sum_{\{i,j\} \subset \mathbb{Z}^d} J_{i,j} \sigma_i \sigma_j$$

with $\sigma := (\sigma_i)_{i \in \mathbb{Z}^d} \in \{\pm 1\}$, $J_{i,j} \geq 0$ and the Boltzmann distribution:

$$\mathbb{P}_\beta(\omega) \propto e^{-\beta H(\omega)}.$$

with $\beta \geq 0$ and some configuration $\omega \in \{\pm 1\}^{\mathbb{Z}^d}$.

We also define the inverse *critical temperature* by

$$\beta_c = \inf\{\beta \geq 0 : \inf_{x \in \mathbb{Z}^d} \langle \sigma_0 \sigma_x \rangle_\beta > 0\}.$$

If there exists R such that If $J_{i,j} = 0$ for $\|i - j\|_\infty \geq R$, we say that the model is finite-range.

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► Any local functions f and g , there exist $c_A^f, c_B^g \in \mathbb{R}$ such that

$$f = \sum_{A \subset \text{supp}(f)} c_A^f \sigma_A \quad g = \sum_{B \subset \text{supp}(g)} c_B^g \sigma_B,$$

with $\sigma_A := \prod_{i \in A} \sigma_i$. In particular, one has

$$\langle f; g \rangle_\beta := \text{Cov}_{\mathbb{P}_\beta}[f, g] = \sum_{\substack{A \subset \text{supp}(f) \\ B \subset \text{supp}(g)}} c_A^f c_B^g \langle \sigma_A; \sigma_B \rangle_\beta.$$

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Therefore, understanding $\langle f; g \rangle_\beta$ **amounts to** understanding $\langle \sigma_A; \sigma_B \rangle_\beta$.

► We will take $\beta \neq \beta_c$. Moreover, we will take J super-exponentially decaying: for any $c > 0$, one has

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Question: What can be said about the asymptotic behaviour of $\langle \sigma_0; \sigma_x \rangle_\beta$ as $\|x\| \rightarrow \infty$?

Two-point function: Ornstein and Zernike

► In 1914 and 1916, Ornstein and Zernike developed a (heuristic) theory of correlations with quickly decaying interactions. In particular, they concluded that, at large distances away from the critical temperature, the spin-spin correlation of the Ising model satisfies

$$\langle \sigma_0; \sigma_x \rangle_\beta \sim \|x\|^{-(d-1)/2} e^{-\nu_\beta(x)}.$$

► **Is it possible** to establish this result **rigorously** ?



- One has the following **Ornstein–Zernike asymptotics**:

Theorem

[A.–OTT–VELENIK 2021]

Assume that $\beta < \beta_c$. Let $\vec{s} \in \mathbb{S}^{d-1}$. Then, as $n \rightarrow \infty$,

$$\langle \sigma_0 ; \sigma_{n\vec{s}} \rangle_\beta = \frac{\Psi_\beta(\vec{s})}{n^{(d-1)/2}} e^{-\nu_\beta(\vec{s})n} (1 + o(1)),$$

where the functions Ψ_β and $\nu_\beta(\vec{s})$ are positive and analytic in $\vec{s}$.

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- The above result has a long history. Some milestones are

- ▷ **Ornstein–Zernike 1914, Zernike 1916**: first (non-rigorous) derivation
- ▷ **Wu 1966, Wu et al 1976**: exact computation, planar model, $\beta < \beta_c$
- ▷ **Abraham–Kunz 1977, Paes–Leme 1978**: any dimension, n.n. model, $\beta \ll 1$
- ▷ **Campanino–Ioffe–Velenik 2003**: any dimension, finite range, $\beta < \beta_c$
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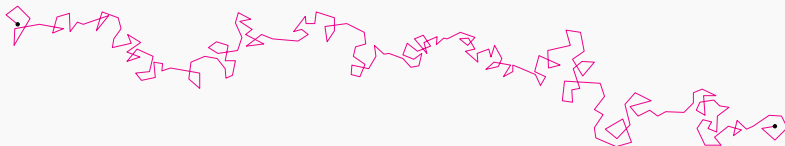
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Probabilistic picture behind OZ asymptotics

- In order to study the subcritical Ising model, one can different graphical representations, for instance the **high temperature expansion**

$$\langle \sigma_0; \sigma_x \rangle_\beta = \sum_{\gamma: 0 \rightarrow x} q_\beta(\gamma).$$

- The techniques developed during last two decades allow to **couple “structurally 1D objects”** with good mixing properties **to directed random walks**.

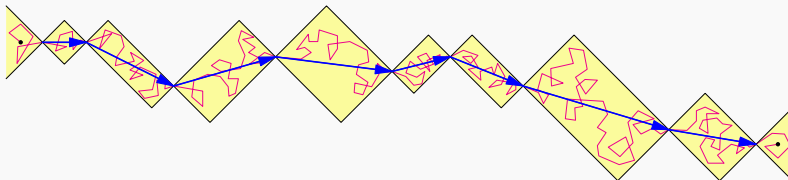


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- The techniques developed during last two decades allow to **couple “structurally 1D objects”** with good mixing properties **to directed random walks**.
- Using this coupling, we can in many cases **reduce difficult questions arising in the Ising model to much simpler (and more classical) ones about random walks**.

► Given $A, B \in \mathbb{Z}^d$ and $\vec{s} \in \mathbb{S}^{d-1}$, we investigate the asymptotic behavior of

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- Of course, by symmetry, $\langle \sigma_C \rangle_\beta = 0$ whenever $|C|$ is odd.

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- We are thus left with **two cases to consider**:

Odd-odd correlations

$|A|, |B|$ both odd

Even-even correlations

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- In the Odd-Odd case, the OZ asymptotics still hold.

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- ▶ The analysis started with the case $|A| = |B| = 2$. Physicists quickly understood that

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- ▶ However, concerning the prefactor, two conflicting predictions were put forward:

Polyakov 1969		Camp, Fisher 1971
n^{-2}	$d = 2$	n^{-d} for all $d \geq 2$
$(n \log n)^{-2}$	$d = 3$	
$n^{-(d-1)}$	$d \geq 4$	

(Note that these predictions only coincide when $d = 2$, where they both agree with the exact computation obtained in *Stephenson 1966* and *Hecht 1967*.)

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- ▶ It turns out that Polyakov was right. This was first shown in
 - ▷ *Bricmont–Fröhlich 1985*: $|A| = |B| = 2$ $\beta \ll 1$ $d \geq 4$
 - ▷ *Minlos–Zhizhina 1988, 1996*: $|A|, |B|$ even $\beta \ll 1$ $d \geq 2$

► The best nonperturbative result to date is the following:

$$\text{Let } \tau(n) = \begin{cases} n^2 & \text{when } d = 2, \\ (n \log n)^2 & \text{when } d = 3, \\ n^{d-1} & \text{when } d \geq 4. \end{cases}$$

Theorem

[OTT-VELENIK 2019]

Let $d \geq 2$ and $\beta < \beta_c$. Let $A, B \in \mathbb{Z}^d$ with $|A|$ and $|B|$ even and let $\vec{s} \in \mathbb{S}^{d-1}$.

Then, there exist constants $0 < C_- \leq C_+ < \infty$ (depending on $A, B, \vec{s}, \beta$) such that, for all n large enough,

$$\frac{C_-}{\tau(n)} e^{-2\nu_\beta(\vec{s})n} \leq \langle \sigma_A ; \sigma_{B+n\vec{s}} \rangle_\beta \leq \frac{C_+}{\tau(n)} e^{-2\nu_\beta(\vec{s})n}.$$

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- ▶ The OZ decay was proved in the ground state of the quantum Ising model in a strong transverse magnetic field (Kennedy 1991). Could it be done more generally with the modern OZ theory ?

THANK YOU AND BUEN APETITO!